Ph.D. Course on *Vorticity, Vortical Flows and Vortex-Induced Vibrations* Technical University of Denmark, Copenhagen, Denmark *vortex.compute.dtu.dk* August 26-30, 2019

## Vorticity Generation 1: Introduction and Key Concepts

#### **Kerry Hourigan**

Fluids Laboratory for Aeronautical and Industrial Research Monash University, Melbourne, Australia Otto Mønsted Guest Professor, Technical University of Denmark





Fluids Laboratory for Aeronautical and Industrial Researc

#### **Course Schedule**

Monday August 26	Tuesday August 27	Wednesday August 28	Thursday August 29	Friday August 30
9.00 – 9.45 Introduction (KH) Lecture 1 (KH) Vorticity Generation I	9.00 – 9.45 Lecture 5 (KH) Swirling Flow Instabilities	9.00 – 9.45 Lecture 9 (KH) Bluff Body Wakes I	9.00 – 9.45 Lecture 12 (KH) Fluid-Structure Interactions I	9.00 – 9.45 Lecture 15 (JS) Aerodynamics of Wind Turbines I
9.45 – 10.00	9.45 – 10.00	9.45 – 10.00	9.45 – 10.00	9.45 – 10.00
Break	Break	Break	Break	Break
10.00 – 10.45	10.00 – 10.45	10.00 – 10.45	10.00 – 10.45	10.00 – 10.45
Lecture 2 (KH)	Lecture 6 (KH)	Lecture 10 (KH)	Lecture 13 (KH)	Lecture 16 (JS)
Vorticity Generation II	Vortex Breakdown	Bluff Body Wakes II	Fluid-Structure Interactions II	Aerodynamics of Wind Turbines II
10.45 – 11.00	10.45 – 11.00	10.45 – 11.00	10.45 – 11.00	10.45 – 11.00
Break	Break	Break	Break	Break
11.00 – 11.45	11.00 – 11.45	11.00 – 11.45	11.00 – 11.45	11.00 – 11.45
Lecture 3	Lecture 7 (MB)	Lecture 11 (KH)	Lecture 14 (KH)	Lecture 17 (KH)
Vortex Filaments (TL)	Vortex Dynamics I	Vehicle Aerodynamics	Fluid-Structure Interactions III	Concluding Lecture
11.45 – 13.00	11.45 – 13.00	11.45 – 13.00	11.45 – 13.00	11.45 - 13.00
Lunch	Lunch	Lunch	Lunch	Lunch
13.00 – 13.45 Lecture 4 Long and Short Wave Instabilities (TL)	13.00 – 13.45 Lecture 8 (MB) Vortex Dynamics II	13.00 – Social programme and dinner	13.00 – 14.00 Project discussions	13.00 – 14.00 Project discussions
13.45 – 14.00 Break	13.45 – 14.00 Break			
14.00 – 14.45 Project presentation	14.00 – 14.45 Project discussions			

#### **General course objectives**

Vorticity is a measure of the spinning motion of a fluid. Understanding the sources, dynamics and structure of vorticity, in particular the interaction of vortices, is of key importance in fluid mechanics. It is the purpose of the course to give the students an overview of recent research in this field.

#### Learning objectives

A student who has met the objectives of the course will be able to:

- understand the mathematical definition and the physical meaning of vorticity
- understand how vorticity is generated and be able to analyze the vorticity flux near boundaries and interfaces
- understand the dynamics of vortex filaments, be able to perform a simple stability analysis
- understand the concept of vortex breakdown and indentify it in various engineering situations
- understand the generation of vortices in the wake of bluff bodies and be able to identify different types of vortex wakes
- perform a basic analysis of the topology of two-dimensional flow, based on dynamical systems theory and experimental or computational data
- understand how vortices interact with solid bodies, how vibrations are induced and how to model this interaction mathematically
- understand the generation of vortices behind wind turbines, and how this is modelled

### **Aims of the Course**

Participants will:

- Gain an introduction to various (but not all!) aspects of vorticity and vortices
- Learn of recent research in vortex dynamics, bluff body flows, flow induced vibrations, wind turbines, vortex breakdown
- Practise some techniques of coding (Python/Anaconda), flow analysis (PIV, Stability) and hands-on simple construction and experiments in anemometry, flow-induced vibration, and turbines.

#### **Lecture Reading Material**

Batchelor, G. K., An Introduction to Fluid Dynamics, Cambridge University Press, 1967.
Holmén, V. Methods for Vortex Identification. Master's Thesis, Lund University, Lund, Sweden, 2012.
Available online: <a href="http://lup.lub.lu.se/student-papers/record/3241710">http://lup.lub.lu.se/student-papers/record/3241710</a>
Lighthill, M. J., Laminar Boundary Layers (ed. L. Rosenhead), Oxford University Press, 1963.
Morton, B. R., The generation and decay of vorticity, Geophysical & Astrophysical Fluid Dynamics, 28:3, 277 - 308, 1984.

#### **Lecture Objectives**

Learn:

The (vector) definition of vorticity cf angular momentum

The usefulness of considering vorticity (and circulation)

Ways of defining vortices

How vorticity is generated at boundaries, transported and its conservation

**Extension to three dimensions** 

#### **Key Questions**

Q1: What is vorticity (and circulation)?

Q2: Why should we worry about vorticity?

Q3: What is a vortex?

Q4: How is vorticity generated?

Q5: How is vorticity transported, and how does it decay?

**Q6: Is vorticity conserved?** 

### **Q1: What is vorticity?**

**Vorticity** is mathematically defined as the curl of the velocity field and is hence a measure of local rotation of the fluid. This definition makes it a vector quantity.

$$\boldsymbol{\omega}(\mathbf{r},t) = \boldsymbol{\nabla} \times \mathbf{v}(\mathbf{r},t)$$

Qualitatively, we usually associate regions of spinning fluid with vorticity, which we term vortices:



Wing tip vorticity (aviation.stackexchange.com

Vorticity in a rotor wake (Siemens)

Boundary layer vorticity Ulrich Rist, *DOI:10.2514/6.2012-84* 

### **Vorticity is local rotation – different to Angular Momentum**

**Vorticity** is a local "rotation", independent of fluid density

$$\boldsymbol{\omega}(\mathbf{r},t) = \boldsymbol{\nabla} \times \mathbf{v}(\mathbf{r},t)$$

Irrotational flow



Steady – no diffusion of vorticity or momentum

Angular momentum uniform,

ω = 0 for r ≠ 0, v = Γ/(2πr), circulation = Γ

en.wikipedia.org/wiki/Vortex

**Angular momentum** is rotation about a fixed origin (non-local) and density dependent



Solid body rotation, no shear,  $\omega = 2\Omega$ ,  $v = \Omega r$ 

#### Circulation

**Circulation**: line integral of tangential velocity around a closed loop or sum of vorticity over an area A (Stokes' theorem)

Vorticity is therefore circulation density

$$\Gamma = \oint_C \mathbf{V} \cdot d\mathbf{l} = \int_A \boldsymbol{\omega} \, \mathrm{d}A$$



#### **Q2: Why should we worry about vorticity?**





www.youtube.com/watch?v=Ee4-xtX9jUQ

#### Flow features are more clearly visualised and understood in terms of vorticity

"Vortices are the sinews and muscles of fluid motions" – Kuchemann (1965)



Velocity and vorticity fields for flow past a cylinder near a free surface (Reichl 2002)

Vorticity field is independent of the Galilean or inertial frame of reference (velocity field is not)

### **Vortices appear in many important flows**



Net vorticity in the flow domain is zero.



#### Wind loads on buildings Vortex induced vibrations

#### Turbulence







## **Aircraft Vortices**



https://www.youtube.com/watch?v=dfY5ZQDzC5s

### **Q3: What is a vortex?**

We qualitatively recognise a vortex as a vorticity structure that retains a degree of coherence with time as it is advected by the flow.

However, due to viscosity, there is no clear boundary to a vortex.

Plotting the vorticity contours shows where vorticity is, but does not show where the boundary of a vortex lies.





#### Methods of identifying a vortex based on the Velocity Gradient tensor

The velocity gradient tensor  $D_{ij}$  can be split into symmetric  $(e_{ij})$  and antisymmetric  $(\omega_{ij})$  tensors:

$$D_{ij} \equiv \frac{\partial U_j}{\partial x_i} \equiv e_{ij} + \omega_{ij} = e_{ij} + \frac{1}{2} \epsilon_{ijk} \omega_k$$

where

Rate of Strain Tensor: 
$$e_{ij} = \frac{1}{2} \left( \frac{\partial U_i}{\partial x_j} + \frac{\partial U_j}{\partial x_i} \right)$$
  
Axisymmetric Rotation Tensor:  $\omega_{ij} = \frac{1}{2} \left( \frac{\partial U_i}{\partial x_j} - \frac{\partial U_j}{\partial x_i} \right)$   
Vorticity Vector:  $\omega_k = \epsilon_{kmn} \frac{\partial U_n}{\partial x_m}$ 

The characteristic equation for  $D_{ij}$  is given by:

$$|D_{ij} - \lambda I| = \lambda^3 + P \lambda^2 + Q \lambda + R = 0$$

where P, Q and R are the three invariants of the velocity gradient tensor.

Using the decomposition into symmetric and anti-symmetric parts these invariants can be expressed as follows:

$$\begin{split} P &= -tr(\bar{D}) \\ Q &= \frac{1}{2}(tr(\bar{D})^2 - tr(\bar{D}^2)) \; = \; \frac{1}{2} ||\bar{\Omega}||^2 - ||\bar{S}||^2 \\ R &= -det(\bar{D}) \end{split}$$

The discriminant of the characteristic equation is:  $(2)^{3}$ 

$$\Delta = \left(\frac{Q}{3}\right)^3 + \left(\frac{R}{2}\right)^2$$

### Q-criterion; $\Delta$ criterion

### Q-criterion

 ${\cal Q}$  represents the local balance between shear rate strain and vorticity magnitude.

Vortices are defined as areas where Q > 0

i.e., where the vorticity magnitude is greater than the magnitude of the rate-of-strain.

### $\Delta$ -criterion

This criterion defines vortices as "regions in which the eigenvalues of  $D_{ij}$  are complex and the streamline pattern is spiralling or closed". For eigenvalues to be complex, the disciminant must be positive:

 $\Delta = \left(\frac{Q}{3}\right)^3 + \left(\frac{R}{2}\right)^2 > 0 \quad \text{(valid for incompressible flows: } P = 0\text{)}.$ 

## $\lambda_2$ criterion

#### $\lambda_2$ criterion

Seeks pressure minimum but removes effects from unsteady straining and viscosity by discarding these terms.

Taking gradient of Navier-Stokes equations:

 $a_{i,j} = -\frac{1}{\rho} p_{ij} + \nu u_{i,jkk}$  where  $a_{i,j}$  is the acceleration gradient and  $p_{ij}$  is symmetric.

Decomposing  $a_{i,j}$  into symmetric and antisymmetric parts give the vorticity transport equation as the antisymmetric part, and the symmetric part:

$$\frac{DS_{ij}}{Dt} - \nu S_{ij,kk} + \Omega_{ik}\Omega_{kj} + S_{ik}S_{kj} = -\frac{1}{\rho}p_{ij}$$

The first 2 terms on LHS represent unsteady irrotational straining and viscous effects, respectively.

Therefore, only  $S^2 + \Omega^2$  is considered to determine if there is a local pressure minimum that entails a vortex.

A vortex is defined as "a connected region with 2 negative eigenvalues of  $S^2 + \Omega^2$ , which is symmetric and therefore has real eigenvalues only. (Jeong & Hussain, 1995).

By ordering the eigenvalues  $\lambda_1 < \lambda_2 < \lambda_3$ , the definition is equivalent to requiring  $\lambda_2 < 0$ .

The isosurfaces are usually visualised for different values of  $-\lambda_2$ .

### **Other Criteria**

There are others such as Swirling Strength: see Methods for Vortex Identification, Vivianne Holmén, 2012)



Figure 2: Comparison between vorticity (in the bottom row) and swirling strength (in the top row) contours at  $U^* = 3.0$  (left column), 7.0 (middle column) and 10.0 (right column) for  $\alpha = 0^{\circ}$ .

#### **Q4: How is vorticity generated?**



## `Early" treatments of the generation and decay of vorticity





- Noted that at almost all points of the boundary there is a non- Devoted a substantial section to the source of vorticity in zero gradient of vorticity along the normal.
- That tangential vorticity must be created at the boundary, in the direction of the surface isobars and at a rate proportional to the tangential pressure gradient.
- Described the boundary as a distributed source of vorticity, similar to a distributed source of heat.

**BATCHELOR - 1967** 



- Identified the free-slip velocity in the inviscid solution for flow past a body set in motion as the effective source of vorticity in real flows.
- Argued that vorticity cannot be destroyed in the interior of a homogeneous fluid, and this appeared to lead to the concept of loss of vorticity by diffusion to boundaries.



# 3 important previous papers since that will be reference (there are others of course – e.g., Wu & Wu's papers)

Morton, B. R., 'The generation and decay of vorticity', Geophysical & Astrophysical Fluid Dynamics, 28:3, 277 -308, 1984.

Rood, E. P., Myths, math, and physics of free-surface vorticity" in "Mechanics USA 1994" edited by AS Kobayashi, Appl Mech Rev vol 47, no 6, part 2, June 1994.

> Lundgren, T. & Koumoutsakos, P., On the generation of vorticity at a free surface, J. Fluid Mech., 382, pp. 351-366, 1999.

# Vorticity equation does not include generation – vorticity is generated only at boundaries/interfaces

The Helmholtz vorticity equation for an incompressible homogeneous fluid

 $\partial \boldsymbol{\omega} / \partial t + (\mathbf{v} \cdot \nabla) \boldsymbol{\omega} = (\boldsymbol{\omega} \cdot \nabla) \mathbf{v} + v \nabla^2 \boldsymbol{\omega},$ 

Processing term

describes the effects of local amplification (or concentration) of vorticity by vortex filament stretching and local turning of filaments Viscous diffusion spread of vorticity due to viscosity

No true generation term How is vorticity generated at boundaries?

## **Vorticity flux at an interface**



# **MORTON:** Generation of vorticity at a solid boundary is an inviscid process

$$\sigma = \frac{d\boldsymbol{u}}{dt} \cdot \hat{\boldsymbol{t}} + \frac{\partial}{\partial x} \frac{p}{\rho} - \boldsymbol{g} \cdot \hat{\boldsymbol{t}}$$

Rate of vorticity generation = (Acceleration of lower boundary) – (Inviscid acceleration of fluid elements due to pressure/body forces)

Morton, B.R (1984) 'The generation and decay of vorticity'



FIGURE 9 The circulation circuit & for generation at a boundary.

Can our integration loop go into the solid? Can we have an inviscid fluid?

## How can we regard a solid as a fluid?

"The most famous version of the experiment was started in 1927 by Professor Thomas Parnell of the University of Queensland in Brisbane, Australia, to demonstrate to students that some substances that appear to be solid are in fact very-high-viscosity fluids"<sup>1</sup>.



Professor John Mainstone (taken in 1990)

Pitch has a viscosity approximately 230 billion (2.3×10<sup>11</sup>) times that of water



<sup>1</sup>en.wikipedia.org/wiki/Pitch\_drop\_experiment

## Superfluids have zero viscosity



Superfluid fountain – frictionless flow forever

www.youtube.com/watch?v=2Z6UJbwxBZI

Morton (1984): Fluid pressure gradient and relative acceleration lead to vorticity production at boundaries/interfaces

$$\llbracket \sigma \rrbracket = \frac{\mathrm{d}\gamma}{\mathrm{d}t} - \frac{\partial}{\partial s} \llbracket \frac{p}{\rho} \rrbracket$$

• Vorticity is generated at boundaries by the relative acceleration of fluid and wall.

• In a homogeneous fluid, vorticity is generated only at boundaries and its generation is instantaneous and the result of inertial forces.

# Wake vorticity generated at boundary by either: fluid pressure gradient



Note that generation is masked by viscous redistribution of the vorticity immediately upon generation.

# ... or boundary acceleration



Vorticity generated due to a differential velocity being created by plate/fluid acceleration

No-slip condition? Infinitely thin layer of fluid accelerated with plate What is the role of viscosity? Redistributes vorticity following generation

## **Q5:** How is vorticity transported, and how does it decay?

## Impulsively started plate



Vorticity/unit length = U<sub>b</sub>

At t = 0, all the vorticity is generated.

t > 0, vorticity diffuses away from the boundary and remains most concentrated at boundary.

Vorticity is neither lost nor gained to solid boundary.

## **Plane Couette flow: two fluids**



- An amount of vorticity U per unit length of plate is generated at t=0.
- This diffuses out into the layer to produce in time two layers, each with uniform vorticity distribution.
- Thereafter the flow is steady and vorticity is neither generated at the plates nor lost to them.
- Due to different µ, at the centre interface, there is a vorticity jump but velocity is continuous (as at other interfaces).

## **Poiseuille flow – no loss of vorticity at boundaries**



- Vorticity is generated continuously by the tangential pressure gradient at the lower boundary, and at the upper boundary at same rate (sense of normal reversed)
- Each diffuses towards the centre plane where the positive and negative fluxes suffer annihilation.
- The circulation per unit length of channel is zero.

# The decay of vorticity

- Vorticity is not lost by diffusion to boundaries other than in circumstances in which vorticity of counter sign is being generated and suffering immediate cross-diffusive annihilation with pre-existing vorticity.
- The gross circulation in an enveloping contour of a bounded region with outer boundary at rest must always be zero.
- Diffusion will, therefore, always lead in time to the total annihilation of all vorticity, through the diffusion of momentum and the corresponding wall stresses.

## **Q6:** Is vorticity conserved?

- Conservation laws play an important role in our understanding of fluid mechanics, and indeed our understanding of the universe. For example, the fundamental equations of fluid mechanics (continuity, Navier-Stokes and energy) equations are usually interpreted as expressions of the conservation of mass, momentum, and energy, respectively.
- Conservation laws often allow the analysis of a system in terms of inlet and outlet properties, as conserved quantities cannot be created or destroyed in the system interior.

## **Conservation of vorticity**

$$\frac{\mathrm{d}}{\mathrm{d}t}\left(\int_{A}\boldsymbol{\omega}\,\mathrm{d}A + \int_{a}^{b}\boldsymbol{\gamma}\,\mathrm{d}s\right) = \oint_{C}\boldsymbol{\nu}\boldsymbol{\nabla}\boldsymbol{\omega}\cdot\hat{\boldsymbol{n}}\,\mathrm{d}s - \left(\left[\left[\frac{p}{\rho}\right]\right]_{b} - \left[\left[\frac{p}{\rho}\right]\right]_{a}\right)$$



# Vorticity is conserved – only need to know what happens at the boundary to know how much vorticity is inside

- Vorticity conservation appears to be a global conservation property, under appropriate conditions, with vorticity generation occurring on a local scale.
- Vorticity generation of one sign must be balanced by creation of an equal magnitude of opposite sign circulation elsewhere in the domain.
- Transient and steady-state behaviour in many flows can then be understood in terms of this conservation principle: vorticity may be generated on a boundary, however the total amount of vorticity generated on all boundaries is often zero.
- The only means by which vorticity generated on a boundary may disappear from the flow is by annihilation due to cross-diffusion with opposite sign vorticity.



## What happens in 3D with an interface between 2 fluids?

Rate of change of total vorticity in control volume V with outer control volume  $\partial V = S_1 \cup S_2$ 

$$\begin{split} \frac{\mathrm{d}\boldsymbol{\Omega}}{\mathrm{d}t} &= \oint_{\partial V} \boldsymbol{\omega} (\boldsymbol{v}^{\mathrm{b}} - \boldsymbol{u}) \cdot \boldsymbol{n} \mathrm{d}S + \oint_{\partial V} (\boldsymbol{\omega} \cdot \boldsymbol{n}) \boldsymbol{u} \mathrm{d}S + \oint_{\partial V} \boldsymbol{\sigma}_{b} \mathrm{d}S - \oint_{\partial I} \left[ \frac{p}{\rho} + \Phi_{g} \right] \mathrm{d}\boldsymbol{s} \\ &- \oint_{\partial I} \frac{1}{2} [\![\boldsymbol{u} \cdot \boldsymbol{u}]\!] \mathrm{d}\boldsymbol{s} - \oint_{\partial I} \boldsymbol{\gamma} (\boldsymbol{v} \cdot \boldsymbol{\hat{b}}) \mathrm{d}\boldsymbol{s} - \oint_{\partial I} (\boldsymbol{\gamma} \cdot \boldsymbol{\hat{b}}) (\boldsymbol{u} \cdot \boldsymbol{\hat{n}}) \boldsymbol{\hat{n}} \mathrm{d}\boldsymbol{s} \end{split}$$

The only source term is the pressure acceleration, which represents the generation of interface circulation due to the inviscid relative acceleration of fluid elements on each side of the interface (extension of Morton's description in 2D).

Viscous tangential acceleration and vortex stretching produces a change in interface circulation, but these are balanced by a decrease of vorticity in the fluid interior.

#### The 9 Commandments of Morton for flow with a solid wall boundary

- 1. that generation results from tangential acceleration of a boundary, from tangential initiation of boundary motion and from tangential pressure gradients acting along a boundary;
- 2. that generation is instantaneous;
- 3. that vorticity once generated cannot subsequently be lost by diffusion to boundaries;
- 4. that reversal of the sense of acceleration or of the sense of pressure gradient results in reversal of the sense of vorticity generated;
- 5. that wall stress relates to the presence of vorticity but is not a cause of its generation;
- 6. that the generation process is independent of the prior presence of vorticity;
- 7. that both senses of vorticity are needed to explain observations;
- 8. that walls play no direct role in the decay or loss of vorticity;
- 9. that vorticity decay results from cross-diffusion of two fluxes of opposite sense and takes place in the fluid interior.

#### Need a tenth Commandment for a free surface boundary

#### References from: Holmén, V. Methods for Vortex Identification. Master's Thesis, Lund University, Lund, Sweden, 2012. Available online: http://lup.lub.lu.se/student-papers/record/3241710

#### References

- Beju, I. (1983) Euclidean tensor calculus with applications. Revised and updated translation of "Tehnici de calcul tensorial euclidian cu aplicatii", Abacus Press 1983
- [2] Pinaki Chakraborty, S. Balachandar, Ronald J. Adrian, 2005 On the relationships between local vortex identification schemes J. Fluid Mech., pp.189-214
- [3] M. S. Chong, A. E. Perry, B. J. Cantwell (1990) A general classification of threedimensional flow fields Phys. Fluids A 1, 765 (1990); doi: 10.1063/1.857730
- [4] R. Cucitore, M. Quadrio, A. Baron, 1999 On the effectiveness and limitations of local criteria forthe identification of a vortex Eur. J. Mech. B/Fluids, 261-282
- [5] Robert L. Devaney An Introduction to Chaotic Dynamical Systems 1986 The Benjamin/Cummings Publishing Company, Inc.
- [6] M.A. Green, C.W Rowley and G. Haller (2007) Detection of Lagrangian coherent structures in three-dimensional turbulence J. Fluid Mech. (2007), vol. 572, pp. 111-120; doi: 10:1017/S0022112006003648
- G. Haller, 2002 Lagrangian coherent structures from approximate velocity data Phys. Fluids 14, 1851 (2002); doi: 10.1063/1.147449
- [8] G. Haller, 2005 An Objective Definition of a Vortex J. Fluid Mech. 1-26
- [9] J.C.R. Hunt, A.A. Wray, P. Moin, 1988 Eddies, stream, and convergence zones in turbulent flows. Center for Turbulence Research Report CTR-S88, pp. 193-208
- [10] M. C. Irwin Smooth Dynamical Systems (2000) First published in 1980. World Scientific Publishing Co.
- [11] Jinhee Jeong, Fazle Hussain, 1995 On the identification of a vortex J. Fluid Mech., pp.69-94
- [12] Jurgen Jost (2005) Dynamical Systems. Examples of Complex Behaviour ( Springer-Verlag Berlin Heidelberg 2005
- [13] S. Kida, H. Miura, 1998 Identification and Analysis of Vortical Structures Eur. J. Mech. B/Fluids, 471-488
- [14] Václav Kolář, 2007 Vortex identification: New requirements and limitations International Journal of Heat and Fluid flow 638-652
- [15] H.J. Lugt, 1979 The dilemma of defining a vortex In Recent Developments in Theoretical and Experimental Fluid Mechanics (ed. U. Mller, K.G. Roesner & B. Schmidt), pp. 309-321. Springer
- [16] Parviz Moin, John Kim, 1984 The Structure of the Vorticity Field in Turbulent Channel Flow Part 1: Analysis of Instantaneous Fields and Statistical Correlations

- [17] Stephen B. Pope (2000) *Turbulent Flows* (1st Edition) Cambridge University Press
- [18] Shawn C. Shadden, Francois Lekien, Jerrold E. Marsden (2005) Definition and properties of Lagrangian coherent structures from finite-time Lyapunov exponents in two-dimensional aperiodic flows Physica D 212 (2005) 271-304
- [19] C. Truesdell, 1953 The Kinematics of Vorticity Indiana University.
- [20] Donald F. Young, Bruce R. Munson, Theodore H. Okiishi, Wade W. Huebsch (2007) A Brief Introduction to Fluid Mechanics (4th Edition) John Wiley & sons Inc.
- [21] J. Zhou, R.J.Adrian, S. Balachandar, T.M. Kendall (1999) Mechanisms for generation coherent packets of hairpin vortices in channel flow Journal of Fluid Mechanics, 387, pp 353-396 doi:10.1017/S002211209900467X